On the contrary, we can prove that a soluble model such as Corngold's[2] is stable against small changes of $\Sigma_s(v)$ eventually due to experimental uncertainties in measuring the cross sections. The model used by Corngold,

$$v'\Sigma_s(v', v) = \beta v'\Sigma_s(v')\, v\Sigma_s(v)\, M(v) \quad ,$$

leads to

$$\tilde{n}_\lambda = \beta M(v)\, \frac{v\Sigma_s(v)v_0\Sigma_s(v_0)}{\lambda + v_0\Sigma(v_0) + iBv_0\mu_0}\, \frac{g(v,\lambda)}{1 - \beta b(\lambda)}$$

$$+ \delta(v - v_0)\, \frac{1}{\lambda + v\Sigma(v) + iBv\mu_0}$$

with

$$g(v,\lambda) = \frac{1}{2iBv} \ln \frac{\lambda + v\Sigma + iBv}{\lambda + v\Sigma - iBv}$$

$$b(\lambda) = \int_0^\infty (v\Sigma_s)^2 M(v)\, g(v,\lambda)\, dv \quad .$$

Let $\Sigma_s'(v)$ differ from $\Sigma_s(v)$ by an error $\epsilon(v)$ which is constant for simplicity: $\Sigma_s'(v) = \Sigma_s(v) + \epsilon$. Then

$$b'(\lambda) = b(\lambda) + \int_0^\infty dv(2v^2 M\Sigma_s\epsilon + v^2\epsilon^2 M)\, g(v,\lambda)$$

$$= b(\lambda) + \eta_\epsilon(\lambda) \quad \lim_{\epsilon \to 0} \eta_\epsilon(\lambda) \to 0$$

and $\delta\tilde{n}_\lambda = \tilde{n}_\lambda' - \tilde{n}_\lambda$ tends to zero when ϵ tends to zero, uniformly in λ.

It would probably be interesting to find the largest class of functions for the sections for which the inverse Laplace transform would be a stable procedure against the small (experimental) errors; however, because we ignore the true scattering kernel in regard to its analytic formula and we have it given at most by a histogram, we do not believe that one can find a function among the elements of this class that could be "brought closer to reality."[3] At most, this class will be a class of mathematical rather than physical models.

V. Protopopescu

Institute for Atomic Physics
Bucharest, Romania

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[2] N. CORNGOLD, *Nucl. Sci. Eng.*, **19**, 80 (1964).

[3] R. CONN and N. CORNGOLD, *Nucl. Sci. Eng.*, **37**, 85 (1969).

Corrigendum

T. C. CHAWLA and B. M. HOGLUND, "Pressure Distribution for Axisymmetric Two-Dimensional Flow in a Plenum During Coolant Explusion," *Nucl. Sci. Eng.*, **43**, 87 (1971).

On p. 89, Eqs. (17a) and (17b) are to be replaced by the following single equation that is applicable for both cases, viz., $m \neq n$ and $m = n$:

$$P_{av}(0,t) = \rho g H + \rho \left[H + \frac{4}{b^2}\sum_{n=1}^{\infty} \frac{J_1^2(\alpha_n b)}{\alpha_n^3 J_0^2(\alpha_n a)}\right] \frac{d}{dt}\left(U_0 \frac{b^2}{a^2}\right)$$

$$- \frac{1}{2}\rho \left\{\frac{8b^2}{a^4} U_0^2 \sum_{n=1}^{\infty} \frac{J_1^2(\alpha_n b)}{\alpha_n^2 J_0^4(\alpha_n a)} \left[J_0^2(\alpha_n b) + J_1^2(\alpha_n b) - \frac{1}{\alpha_n b} J_0(\alpha_n b)\, J_1(\alpha_n b)\right]\right\}$$

$$- \frac{1}{2}\rho \left\{\frac{16b}{a^4}\, U_0^2 \sum_{m=1}^{\infty} \sum_{n=m+1}^{\infty} \frac{J_1(\alpha_m b)\, J_1(\alpha_n b)}{\alpha_m \alpha_n(\alpha_m - \alpha_n)\, J_0^2(\alpha_m a)\, J_0^2(\alpha_n a)}\right.$$

$$\times [J_1(\alpha_m b)\, J_0(\alpha_n b) - J_0(\alpha_m b)\, J_1(\alpha_n b)]$$

$$\left. + \frac{8b^2}{a^4} U_0^2 \sum_{n=1}^{\infty} \frac{J_1^2(\alpha_n b)}{\alpha_n^2 J_0^2(\alpha_n a)}\right\} + P_a \quad . \tag{17}$$